

# Earth and Space Science



## RESEARCH ARTICLE

10.1029/2021EA001763

### Key Points:

- Paper develops accuracy equations for CO<sub>2</sub>/H<sub>2</sub>O measurements from infrared gas analyzers in closed-path eddy-covariance systems
- The equations are used to assess the improvement in CO<sub>2</sub>/H<sub>2</sub>O measurement accuracy of analyzers under automatic zero/span procedures
- Assessments find that, for analyzers in cold/dry conditions, the zero procedure is important, but the H<sub>2</sub>O span procedure is unnecessary

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### Citation:

Zhou, X., Gao, T., Pang, Y., Mahan, H., Li, X., Zheng, N., et al. (2021). Based on atmospheric physics and ecological principle to assess the accuracies of field CO<sub>2</sub>/H<sub>2</sub>O measurements from infrared gas analyzers in closed-path eddy-covariance systems. *Earth and Space Science*, 8, e2021EA001763. <https://doi.org/10.1029/2021EA001763>

Received 24 MAR 2021  
 Accepted 21 SEP 2021

## Based on Atmospheric Physics and Ecological Principle to Assess the Accuracies of Field CO<sub>2</sub>/H<sub>2</sub>O Measurements From Infrared Gas Analyzers in Closed-Path Eddy-Covariance Systems

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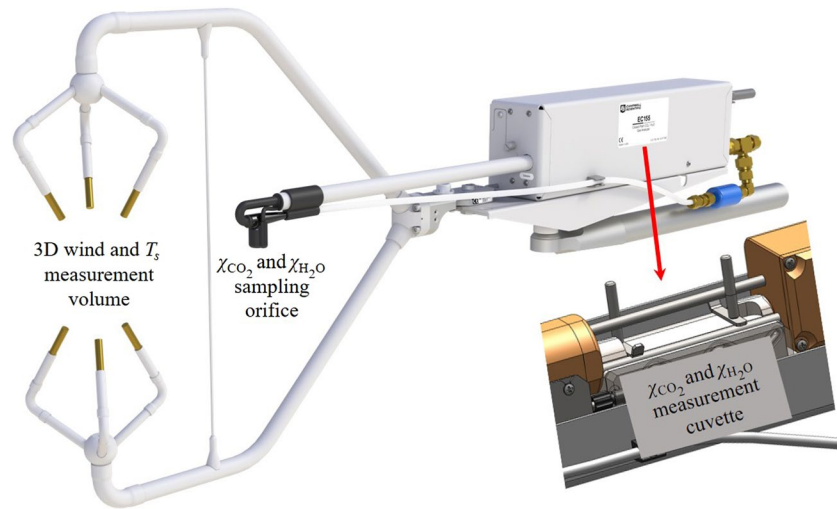
**Abstract** Field CO<sub>2</sub>/H<sub>2</sub>O measurements from infrared gas analyzers in closed-path eddy-covariance systems have wide applications in earth sciences. Knowledge about exactness of these measurements is required to assess measurement applicability. Although the analyzers are specified with uncertainty components (zero drift, gain drift, cross-sensitivities, and precision), exactness for individual measurements is unavailable due to an absence of methodology to comprehend the components as an overall uncertainty. Adopting an advanced definition of accuracy as a range of all measurement uncertainty sources, the specified components are composited into a model formulated for studying analyzers' CO<sub>2</sub>/H<sub>2</sub>O accuracy equations. Based on atmospheric physics and environmental parameters, the analyzers are evaluated using the equations for CO<sub>2</sub> accuracy ( $\pm 0.78 \mu\text{molCO}_2 \text{ mol}^{-1}$ , relatively  $\pm 0.18\%$ ) and H<sub>2</sub>O accuracy ( $\pm 0.15 \text{ mmolH}_2\text{O mol}^{-1}$ ). Evaluation shows that precision and cross-sensitivity are minor uncertainties while zero and gain drifts are major uncertainties. Both drifts need adjusting through zero/span procedures during field maintenance. The equations provide rationales to guide and assess the procedures. H<sub>2</sub>O span needs more attentions under humid conditions. Under freezing conditions while H<sub>2</sub>O span is impractical, this span is fortunately unnecessary. Under the same conditions, H<sub>2</sub>O zero drift dominates H<sub>2</sub>O measurement uncertainty. Therefore, automatic zero becomes a more applicable and necessary tactic. In general cases of atmospheric CO<sub>2</sub> background, automatic CO<sub>2</sub> zero/span procedures can narrow CO<sub>2</sub> accuracy by 36% ( $\pm 0.74$  to  $\pm 0.47 \mu\text{molCO}_2 \text{ mol}^{-1}$ ). Automatic/manual H<sub>2</sub>O zero/span procedures can narrow H<sub>2</sub>O accuracy by 27% ( $\pm 0.15$  to  $\pm 0.11 \text{ mmolH}_2\text{O mol}^{-1}$ ). While ensuring system specifications, the procedures guided by equations improve measurement accuracies.

## 1. Introduction

Closed-path eddy-covariance (CPEC) systems are used to measure boundary-layer CO<sub>2</sub>, H<sub>2</sub>O, heat, and momentum fluxes between ecosystems and the atmosphere (Ibrom et al., 2007; Leuning & Moncrieff, 1990). For the fluxes, a CPEC system is equipped with a fast response three-dimensional (3D) sonic anemometer to measure wind and sonic temperature ( $T_s$ ) and a fast-response infrared gas analyzer to measure CO<sub>2</sub> and H<sub>2</sub>O amounts. In this configuration, CO<sub>2</sub> and H<sub>2</sub>O are measured inside the analyzer cuvette. For both measurements, air is sampled into the cuvette from the analyzer sampling orifice adjacently positioned to the sonic measurement volume (Figure 1). Together, the anemometer and analyzer provide high-frequency (e.g., 10 Hz) measurements used to compute the fluxes (Aubinet et al., 2012) at a location represented by the sonic measurement volume and the gas analyzer orifice position. The degree of exactness of each flux from the measured data depends primarily on the exactness of field measurements for CO<sub>2</sub>, H<sub>2</sub>O, and/or  $T_s$  along with 3D wind (Fratini et al., 2014). Beyond the acquisition for the fluxes, the data of individual variables from these field measurements have various applications in other domains. In many settings, knowledge of measurement exactness is required for assessing data applicability. This study models and assesses this exactness of CO<sub>2</sub>/H<sub>2</sub>O data from the infrared gas analyzers in CPEC systems (Figure 1) used in ecosystems.

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**Figure 1.** Sonic measurement volume for three-dimensional (3D) wind and sonic temperature ( $T_s$ ), gas analyzer sampling orifice position for  $\text{CO}_2$  mixing ratio ( $\chi_{\text{CO}_2}$ ) and  $\text{H}_2\text{O}$  mixing ratio ( $\chi_{\text{H}_2\text{O}}$ ), and the gas analyzer measurement cuvette for  $\chi_{\text{CO}_2}$  and  $\chi_{\text{H}_2\text{O}}$  in a CPEC system (e.g., CPEC310, Campbell Scientific Inc., UT, USA).

The analyzers in CPEC systems output  $\text{CO}_2$  mixing ratio (i.e.,  $\chi_{\text{CO}_2} = \rho_{\text{CO}_2} / \rho_d$  where  $\rho_{\text{CO}_2}$  is  $\text{CO}_2$  molar concentration and  $\rho_d$  is dry air molar concentration) and  $\text{H}_2\text{O}$  mixing ratio (i.e.,  $\chi_{\text{H}_2\text{O}} = \rho_{\text{H}_2\text{O}} / \rho_d$ , where  $\rho_{\text{H}_2\text{O}}$  is  $\text{H}_2\text{O}$  molar concentration). For instance,  $\chi_{\text{H}_2\text{O}}$  along with  $T_s$  can be used to derive ambient air temperature ( $T_a$ ) (Kaimal & Gaynor, 1991; Schotanus et al., 1983). In this case, given an exact  $T_a$  equation, the applicability of equation relies solely on the measurement exactness of  $\chi_{\text{H}_2\text{O}}$  and  $T_s$ . Therefore, the higher the degree of exactness of  $\chi_{\text{H}_2\text{O}}$  and  $T_s$ , the higher the degree of exactness of  $T_a$ . The evaluation on the uncertainty of  $T_a$  from  $\chi_{\text{H}_2\text{O}}$  and  $T_s$  measurements needs the overall exactness values of  $\chi_{\text{H}_2\text{O}}$  and  $T_s$ . Although the uncertainty sources related to the exactness of gas analyzer measurements are separately specified by drifts, cross-sensitivities, and precision (Campbell Scientific Inc., 2018a; LI-COR Biosciences, 2016), specifications for the exactness of such individual field measurements have been unavailable until now. This is due to the absence of methodology to comprehend all individual uncertainty sources.

For any sensor, the measurement exactness depends on its performances, which are commonly specified in terms of accuracy, precision, and other uncertainty descriptors, such as sensor drift. Conveniently, the accuracy represents trueness as a systematic uncertainty to quantify the degree of closeness of measurement to the true value in a measured quantity. The precision is the standard deviation of random measurement errors to quantify the degree to which repeated measurements under unchanged conditions show the same results (Joint Committee for Guide in Metrology, 2008). Both accuracy and precision are universally applicable to any sensor for its performance specifications. Other uncertainty descriptors are more sensor specific. For example, cross-sensitivity of  $\text{CO}_2$  measurement detection to  $\text{H}_2\text{O}$  may be applicable only to infrared gas analyzers. Overall exactness of individual measurements is comprehensively descriptive and, in practice as in the above case for  $T_a$ , is most needed by users for their data analyses. Therefore, International Organization for Standardization (2012) advanced the definition of accuracy in a comprehensive way. The accuracy was expanded in its definition as the combination of both trueness and precision. This advanced definition is advantageous and, while keeping the terminology of “accuracy,” consolidates all measurement uncertainties together. Adopting this definition, we specify the accuracy of individual measurements from infrared gas analyzers as the range of total uncertainty from all individual uncertainty sources in field measurements. Using the analyzer specifications of the CPEC300 series (Campbell Scientific Inc.) as an example, we: (a) develop methodologies to comprehend all measurement uncertainty sources of infrared gas analyzers as the accuracy of  $\text{CO}_2/\text{H}_2\text{O}$  measurements into an equation, (b) assess the accuracy of  $\text{CO}_2/\text{H}_2\text{O}$  measurements using the equation, and (c) discuss applications of the assessments in data analyses and analyzer field maintenance. Additional objective of this study is to find an approach for flux community to assess the accuracy of field  $\text{CO}_2/\text{H}_2\text{O}$  measurements from infrared gas analyzers in CPEC systems.

**Table 1**  
Measurement Specifications for EC155 Infrared CO<sub>2</sub>/H<sub>2</sub>O Analyzers (Campbell Scientific Inc., 2018a)

	CO <sub>2</sub>			H <sub>2</sub> O			Note
	Notation	Value	Unit	Notation	Value	Unit	
Precision	$\sigma_{CO_2}$	0.15	$\mu\text{molCO}_2 \text{ mol}^{-1}$ <sup>a</sup>	$\sigma_{H_2O}$	$6.0 \times 10^{-3}$	$\text{mmolH}_2\text{O mol}^{-1}$	
Zero drift	$d_{cz}$	$\pm 0.30$	$\mu\text{molCO}_2 \text{ mol}^{-1}$	$d_{wz}$	$\pm 0.05$	$\text{mmolH}_2\text{O mol}^{-1}$	Both zero and gain drift values are the possible maxima within the system operational range of ambient air temperature. The actual values depend more on this temperature.
Gain drift	$d_{cg}$	$\pm 0.10\%$ <sup>b</sup> true $\chi_{CO_2}$	$\mu\text{molCO}_2 \text{ mol}^{-1}$	$d_{wg}$	$\pm 0.30\%$ <sup>c</sup> true $\chi_{H_2O}$	$\text{mmolH}_2\text{O mol}^{-1}$	
Cross-sensitivity to H <sub>2</sub> O	$s_w$	$\pm 5.6 \times 10^{-8}$	$\mu\text{molCO}_2 \text{ mol}^{-1} (\text{mmolH}_2\text{O mol}^{-1})^{-1}$		N/A		
Cross-sensitivity to CO <sub>2</sub>		N/A		$s_c$	$\pm 5.0 \times 10^{-5}$	$\text{mmolH}_2\text{O mol}^{-1} (\mu\text{molCO}_2 \text{ mol}^{-1})^{-1}$	
Calibration range		0–1,000	$\mu\text{molCO}_2 \text{ mol}^{-1}$		0–79	$\text{mmolH}_2\text{O mol}^{-1}$	For CO <sub>2</sub> , up to 3,000 $\mu\text{molCO}_2 \text{ mol}^{-1}$ if specially needed.

<sup>a</sup>mol in the denominator of all units is mole of dry air. <sup>b</sup>0.10% is CO<sub>2</sub> gain drift percentage denoted by  $\delta_{CO_2-g}$  in text and  $\chi_{CO_2}$  is CO<sub>2</sub> molar mixing ratio. <sup>c</sup>0.30% is H<sub>2</sub>O gain drift percentage denoted by  $\delta_{H_2O-g}$  in text and  $\chi_{H_2O}$  is H<sub>2</sub>O molar mixing ratio.

## 2. Specifications

A system of CPEC300 series includes, but is not limited to, a CSAT3A 3D sonic anemometer and an EC155 infrared CO<sub>2</sub>/H<sub>2</sub>O analyzer. The system operates in an ambient air temperature range of –30 to 50°C and an atmospheric pressure range of 70–106 kPa. The specifications for CO<sub>2</sub> and H<sub>2</sub>O measurements are given in Table 1.

In this table, the top limit of 1,000  $\mu\text{molCO}_2 \text{ mol}^{-1}$  in the calibration range for CO<sub>2</sub> is more than double the background CO<sub>2</sub> mixing ratio in the atmosphere (415  $\mu\text{molCO}_2 \text{ mol}^{-1}$ , Global Monitoring Laboratory, 2021). The top limit of 79  $\text{mmolH}_2\text{O mol}^{-1}$  in the calibration range for H<sub>2</sub>O is equivalent to 40°C dew point under the atmospheric pressure of 101.325 kPa as used by Wright et al. (2003). This limit is higher than the highest dew point of 35°C ever recorded under natural conditions on Earth (National Weather Service, 2021).

The uncertainties of gas analyzers for CO<sub>2</sub> and H<sub>2</sub>O measurements in Table 1 are specified by individual uncertainty sources along with their magnitudes. For CO<sub>2</sub> and H<sub>2</sub>O measurements, respectively, the composite uncertainty range (i.e., the accuracy) is needed most and should be derived from these sources.

The precision uncertainty is caused by random measurement errors, and the other uncertainties can be considered as systematic uncertainties related to trueness. As noted in Table 1, zero and gain drifts are more influenced by ambient air temperature. Additionally, each gain drift is also positively proportional to its own magnitude (i.e., true  $\chi_{CO_2}$  or  $\chi_{H_2O}$ ). Lastly, while measuring CO<sub>2</sub>, sensitivity-to-H<sub>2</sub>O is related to the background concentration of H<sub>2</sub>O as indicated by its unit of  $\mu\text{molCO}_2 \text{ mol}^{-1} (\text{mmolH}_2\text{O mol}^{-1})^{-1}$  and, while measuring H<sub>2</sub>O, sensitivity-to-CO<sub>2</sub> is related to the background concentration of CO<sub>2</sub> as indicated by its unit of  $\text{mmolH}_2\text{O mol}^{-1} (\mu\text{molCO}_2 \text{ mol}^{-1})^{-1}$ .

Accordingly, beyond statistical analysis, the accuracy of CO<sub>2</sub>/H<sub>2</sub>O measurements should be evaluated in an ambient air temperature range of –30 to 50°C, a  $\chi_{H_2O}$  range of 0–79  $\text{mmolH}_2\text{O mol}^{-1}$ , and a  $\chi_{CO_2}$  range up to 1,000  $\mu\text{molCO}_2 \text{ mol}^{-1}$ .

### 3. Accuracy Model

As a maximum range of composite uncertainty, the accuracy is determined collectively by all individual uncertainty components: zero drift, gain drift, precision, and cross-sensitivity-to- $\alpha$  uncertainties where  $\alpha$  can be either CO<sub>2</sub>, if H<sub>2</sub>O is measured, or H<sub>2</sub>O, if CO<sub>2</sub> is measured. Given the true  $\alpha$  mixing ratio ( $\chi_{\alpha T}$ ) and measured one ( $\chi_{\alpha}$ ), the composite uncertainty of measured  $\alpha$  mixing ratios ( $\Delta\chi_{\alpha}$ ) is given by:

$$\Delta\chi_{\alpha} = \chi_{\alpha} - \chi_{\alpha T} \quad (1)$$

The accuracy model is an expression of the maximum range of  $\Delta\chi_{\alpha}$  in terms of quantified measurement uncertainties.

The zero drift uncertainty ( $\Delta\chi_{\alpha}^z$ ) is independent of  $\chi_{\alpha T}$  in value. The cross-sensitivity uncertainty ( $\Delta\chi_{\alpha}^s$ ) is also independent of  $\chi_{\alpha T}$  in value, but depends on the amount of background H<sub>2</sub>O in the air if  $\alpha$  is CO<sub>2</sub> and on the amount of background CO<sub>2</sub> in air if  $\alpha$  is H<sub>2</sub>O. Therefore, while both gain drift and precision uncertainties are zero,  $\Delta\chi_{\alpha}^z$  and  $\Delta\chi_{\alpha}^s$  are additive to  $\chi_{\alpha T}$  as a measured value with zero drift and cross-sensitivity uncertainties ( $\chi_{\alpha-zs}$ ):

$$\chi_{\alpha-zs} = \chi_{\alpha T} + \Delta\chi_{\alpha}^z + \Delta\chi_{\alpha}^s \quad (2)$$

Along with the zero drift and the cross-sensitivity uncertainty in a measurement process, if gain also drifts,  $\chi_{\alpha-zs}$  would be a base magnitude from which gain drifts. As such, the measured value with the zero drift and cross-sensitivity uncertainties plus the gain drift uncertainty ( $\chi_{\alpha-zsg}$ ) can be evaluated as:

$$\chi_{\alpha-zsg} = \chi_{\alpha-zs} + \delta_{\alpha-g}\chi_{\alpha-zs} \quad (3)$$

where  $\delta_{\alpha-g}$  is the gain drift percentage ( $\delta_{CO_2-g} = 0.10\%$  and  $\delta_{H_2O-g} = 0.30\%$ , Table 1). Substituting  $\chi_{\alpha-zs}$  in this equation with Equation 2 leads to:

$$\chi_{\alpha-zsg} = \chi_{\alpha T} + \Delta\chi_{\alpha}^z + \Delta\chi_{\alpha}^s + \delta_{\alpha-g}\chi_{\alpha T} + \delta_{\alpha-g}\Delta\chi_{\alpha}^z + \delta_{\alpha-g}\Delta\chi_{\alpha}^s \quad (4)$$

In the right side of this equation, the magnitude of  $\delta_{\alpha-g}\Delta\chi_{\alpha}^z$  is three orders smaller than  $\Delta\chi_{\alpha}^z$  and the magnitude of  $\delta_{\alpha-g}\Delta\chi_{\alpha}^s$  is three orders smaller than  $\Delta\chi_{\alpha}^s$ . These two smaller terms can be dropped and Equation 4 can be approximated and re-arranged as:

$$\begin{aligned} \chi_{\alpha-zsg} &\approx \chi_{\alpha T} + \Delta\chi_{\alpha}^z + \delta_{\alpha-g}\chi_{\alpha T} + \Delta\chi_{\alpha}^s \\ &= \chi_{\alpha T} + \Delta\chi_{\alpha}^z + \Delta\chi_{\alpha}^g + \Delta\chi_{\alpha}^s \end{aligned} \quad (5)$$

Any measured value has a random error (i.e., precision uncertainty) independent of  $\chi_{\alpha T}$  in value that is caused by unknown minor factors (International Organization for Standardization, 2012). Therefore, the precision uncertainty is additive to any measurement. Adding this precision uncertainty ( $\Delta\chi_{\alpha}^p$ ) to  $\chi_{\alpha-zsg}$  leads to a measured value ( $\chi_{\alpha}$ ) including all uncertainties, given by:

$$\begin{aligned} \chi_{\alpha} &= \chi_{\alpha-zsg} + \Delta\chi_{\alpha}^p \\ &= \chi_{\alpha T} + \Delta\chi_{\alpha}^z + \Delta\chi_{\alpha}^g + \Delta\chi_{\alpha}^s + \Delta\chi_{\alpha}^p \end{aligned} \quad (6)$$

The replacement of  $\chi_{\alpha}$  in Equation 1 with this equation expresses the composite uncertainty as:

$$\Delta\chi_{\alpha} = \Delta\chi_{\alpha}^z + \Delta\chi_{\alpha}^g + \Delta\chi_{\alpha}^s + \Delta\chi_{\alpha}^p \quad (7)$$

$$\Delta\chi_{\alpha} \leq \left| \Delta\chi_{\alpha}^z \right| + \left| \Delta\chi_{\alpha}^g \right| + \left| \Delta\chi_{\alpha}^s \right| + \left| \Delta\chi_{\alpha}^p \right| \quad (8)$$

The four terms in the right side of this equation define a range of composite uncertainty for  $\alpha$  gas species measurements as the accuracy in a model:

$$\Delta\chi_{\alpha} \equiv \pm \left( \left| \Delta\chi_{\alpha}^z \right| + \left| \Delta\chi_{\alpha}^g \right| + \left| \Delta\chi_{\alpha}^s \right| + \left| \Delta\chi_{\alpha}^p \right| \right) \quad (9)$$

Assessment on the accuracy of field CO<sub>2</sub> or H<sub>2</sub>O measurements is to formulate and evaluate the four terms in the right side of this model. The involvement of dry air molar concentration in the expression of  $\chi_{CO_2}$

and  $\chi_{H_2O}$  requires  $H_2O$  molar concentration in moist air to be known first; therefore, the accuracy of  $H_2O$  measurements is studied prior to  $CO_2$ .

#### 4. Accuracy of $H_2O$ Mixing Ratio Measurements

Accuracy model (9) defines the accuracy of  $H_2O$  measurements by infrared gas analyzers ( $\Delta\chi_{H_2O}$ ) as

$$\Delta\chi_{H_2O} \equiv \pm \left( \left| \Delta\chi_{H_2O}^z \right| + \left| \Delta\chi_{H_2O}^g \right| + \left| \Delta\chi_{H_2O}^s \right| + \left| \Delta\chi_{H_2O}^p \right| \right) \quad (10)$$

where  $\Delta\chi_{H_2O}^z$  is  $H_2O$  zero drift uncertainty,  $\Delta\chi_{H_2O}^g$  is  $H_2O$  gain drift uncertainty, and  $\Delta\chi_{H_2O}^s$  is cross-sensitivity-to- $CO_2$  uncertainty, and  $\Delta\chi_{H_2O}^p$  is  $H_2O$  precision uncertainty.

The  $H_2O$  precision is the standard deviation of  $\chi_{H_2O}$  random errors among repeated measurements under the same conditions (International Organization for Standardization, 2012). Accordingly, the precision uncertainty in an individual  $H_2O$  measurement due to this deviation at a  $P$ -value of 0.05 can be defined by statistic theory (Snedecor & Cochran, 1989) as:

$$\Delta\chi_{H_2O}^p = \pm 1.96 \times \sigma_{H_2O} \quad (11)$$

The remaining uncertainties due to  $H_2O$  zero drift,  $H_2O$  gain drift, and cross-sensitivity-to- $CO_2$  are caused by the inability of the working equation inside the gas analyzers to perform consistently for a long-term (e.g., months or seasons) under varying environmental conditions such as, ambient air temperature. According to LI-Cor Biosciences (2016), a general model of the working equation for  $\chi_{H_2O}$  is given by:

$$\begin{aligned} \rho_{H_2O} &= P \sum_{i=1}^3 a_{wi} \left\{ 1 - \left[ \frac{A_w}{A_{ws}} + S_c \left( 1 - \frac{A_c}{A_{cs}} \right) \right] Z_w \right\}^i \left\{ \frac{G_w}{P} \right\}^i \\ \chi_{H_2O} &= \rho_{H_2O} \left[ \frac{P}{R^* (T_g + 273.15)} - \frac{\rho_{H_2O}}{1000} \right]^{-1} \end{aligned} \quad (12)$$

where  $\rho_{H_2O}$  is  $H_2O$  molar concentration in  $mmolH_2O \text{ m}^{-3}$ ;  $a_{wi}$  ( $i = 1, 2, \text{ or } 3$ ) is a coefficient of the third order polynomial in the terms inside curly brackets;  $A_{ws}$  and  $A_{cs}$  are the power of analyzer source lights in the wavelengths for  $H_2O$  and  $CO_2$  measurements, respectively;  $A_w$  and  $A_c$  are the portions of source light power of  $A_{ws}$  and  $A_{cs}$  that pass through the gas;  $S_c$  is cross-sensitivity of detector to  $CO_2$ , while detecting  $H_2O$ , in the wavelength for  $H_2O$  measurements (hereinafter referred as sensitivity-to- $CO_2$ );  $Z_w$  is  $H_2O$  zero adjustment ( $H_2O$  zero coefficient);  $G_w$  is  $H_2O$  gain adjustment (i.e.,  $H_2O$  gain coefficient commonly as  $H_2O$  span coefficient);  $P$  and  $T_g$  are gas pressure and gas temperature, respectively, inside the closed-cuvette; and  $R^*$  is the universal gas constant. The parameters of  $a_{wi}$ ,  $Z_w$ ,  $G_w$ , and  $S_c$  in this model are statistically estimated to establish a  $H_2O$  working equation in the production calibration against a series of standard gases in a range of  $H_2O$  along with  $CO_2$  molar concentrations under a range of  $P$  (hereinafter referred as calibration). The  $H_2O$  working equation (i.e., Model [12] with estimated parameters) is used inside the gas analyzer to compute  $\rho_{H_2O}$  and  $\chi_{H_2O}$  from field measurements of  $A_c$ ,  $A_{cs}$ ,  $A_w$ ,  $A_{ws}$ ,  $P$ , and  $T_g$ .

The working equation is analyzer-specific and is deemed accurate immediately after the calibration process (LI-Cor Biosciences, 2016). However, similar to all optical instruments, after being used in environments different from the manufacturer calibration conditions, an analyzer drifts in  $H_2O$  zero and/or gain. As Model (12) for  $\rho_{H_2O}$  shows, parameter  $Z_w$  is related to  $H_2O$  zero drift;  $G_w$ , to  $H_2O$  gain drift; and  $S_c$ , to sensitivity-to- $CO_2$ . Therefore, the analyses of  $Z_w$  and  $G_w$  along with  $S_c$  are an approach to understand the causes of zero drift, gain drift, and sensitivity-to- $CO_2$ . Such understanding is essential to formulate  $\Delta\chi_{H_2O}^z$ ,  $\Delta\chi_{H_2O}^g$ , and  $\Delta\chi_{H_2O}^s$ .

##### 4.1. $Z_w$ and $\Delta\chi_{H_2O}^z$ ( $H_2O$ Zero Drift Uncertainty)

Gas analyzers are calibrated to report zero  $\chi_{H_2O}$  plus the precision uncertainty for zero gas that is free of  $H_2O$  and  $CO_2$  (hereinafter referred as zero gas). However, when used in measurement conditions that are vastly different from the calibration conditions, the analyzers often report non-zero  $\chi_{H_2O}$  value for zero gas, even

beyond  $\pm\Delta\chi_{H_2O}^p$ . This instability of gas analyzers is termed as H<sub>2</sub>O zero drift. The drift is primarily affected by air temperature surrounding the gas analyzer that is different from the ambient air temperature in the calibration processes ( $T_c$ ) and/or by small H<sub>2</sub>O accumulation inside the analyzer light housing (hereinafter referred as housing H<sub>2</sub>O accumulation) due to unavoidable little gas leaking during long-term use. The light housing is technically sealed to keep housing air close to zero gas by using molecular sieve to remove CO<sub>2</sub> and H<sub>2</sub>O from any ambient air that may sneak into the housing (Campbell Scientific Inc., 2018a).

Due to the H<sub>2</sub>O zero drift, the working equation needs to be adjusted through its parameter re-estimation to adapt to the ambient air temperature and housing H<sub>2</sub>O accumulation near which the system is running. This adjustment is the zero procedure to bring  $\chi_{H_2O}$  and  $\chi_{CO_2}$  of the zero gas from the working equation back to zero, or as close as possible. This section just focuses on H<sub>2</sub>O instead of CO<sub>2</sub> for our discussion about the zero procedure. The same theory applies for CO<sub>2</sub>.

In the field, a simple zero procedure is preferred. Since only a zero H<sub>2</sub>O value is available, the simplest method is to use zero gas to re-estimate one parameter in the working equation that results in zero  $\chi_{H_2O}$  due to zero  $\rho_{H_2O}$ . As Model (12) shows, this parameter for H<sub>2</sub>O turns out to be  $Z_w$  adjustable to result in zero  $\rho_{H_2O}$  for zero gas if re-estimated by:

$$Z_w = \left[ \frac{A_{w0}}{A_{ws}} + S_c \left( 1 - \frac{A_{c0}}{A_{cs}} \right) \right]^{-1} \quad (13)$$

where  $A_{w0}$  and  $A_{c0}$  are the counterparts of  $A_w$  and  $A_c$  for zero gas, respectively. Inside the analyzer, the zero procedure for H<sub>2</sub>O is to re-estimate the H<sub>2</sub>O zero coefficient to satisfy Equation 13.

If the H<sub>2</sub>O zero coefficient always satisfies Equation 13 after the zero procedure, the H<sub>2</sub>O zero drift would not cause a significant uncertainty in H<sub>2</sub>O measurements; however, this is not the case. Similar to the performance after the calibration, an analyzer after the zero procedure will likely drift slowly under changing ambient air temperature. Nevertheless, the value of H<sub>2</sub>O zero coefficient that should be used with the ambient air temperature surrounding the gas analyzer, and particularly with housing H<sub>2</sub>O accumulation, is unpredictable. Given that the molecular sieve inside the analyzer light housing is replaced as recommended in the analyzer maintenance schedule, the housing H<sub>2</sub>O accumulation should not be a concern while the temperature surrounding the gas analyzer is not under control. Therefore, the H<sub>2</sub>O zero drift uncertainty is specified as the maximum range of H<sub>2</sub>O zero drift for the analyzers ( $d_{wz}$ ) that varies with ambient air temperature (Campbell Scientific Inc., 2018b), but normally within the specified range.

Given that an analyzer performs best, almost without zero drift, at the same ambient air temperature as the calibration/zeroing ambient air temperature ( $T_c$ ) and possibly drifts while  $T_g$  changes away from  $T_c$ . The further  $T_g$  is away from  $T_c$ , the more likely it will drift in proportion to the difference between  $T_g$  and  $T_c$  but within the specification over the analyzer operational range of ambient air temperature. Accordingly, H<sub>2</sub>O zero drift uncertainty can be approximated for its maximum range as:

$$\Delta\chi_{H_2O}^z \equiv d_{wz} \times \begin{cases} \frac{T_g - T_c}{T_{rh} - T_{rl}} & T_c < T_g < T_{rh} \\ \frac{T_c - T_g}{T_{rh} - T_{rl}} & T_c > T_g > T_{rl} \end{cases} \quad (14)$$

where, over the analyzer operational range of ambient air temperature,  $T_{rh}$  is the high-end value (50°C for our study case) and  $T_{rl}$  is the low-end value (−30°C for our study case). In this equation,  $\Delta\chi_{H_2O}^z \leq d_{wz}$  over the full range of ambient air temperature from  $T_{rl}$  to  $T_{rh}$  and  $\Delta\chi_{H_2O}^z = d_{wz}$  if  $T_g$  and  $T_c$  are at the two ends of the range (i.e.,  $T_{rl}$  and  $T_{rh}$ ), respectively.

#### 4.2. $G_w$ and $\Delta\chi_{H_2O}^g$ (H<sub>2</sub>O Gain Drift Uncertainty)

All CO<sub>2</sub>/H<sub>2</sub>O analyzers are calibrated against a series of moist air with known H<sub>2</sub>O molar concentrations at different levels. This calibration sets the working equation to closely follow the gain trend in H<sub>2</sub>O change of measured moist air. Similar to the zero drift, during use with changing ambient conditions, the reported gain trend of  $\chi_{H_2O}$  for H<sub>2</sub>O changes in air will possibly drift away from the real gain trend of the change,

which is specifically termed as H<sub>2</sub>O gain drift. This drift is affected by almost the same factors as the H<sub>2</sub>O zero drift (LI-COR Bioscience, 2016).

Due to possible gain drift, the gas analyzer after the zero procedure needs to be further adjusted to tune its working equation back to the real gain trend in H<sub>2</sub>O of measured air or as close as possible, which is the H<sub>2</sub>O span procedure. Like the zero procedure, this procedure is also required to be simple using one H<sub>2</sub>O span gas with known water density ( $\tilde{\rho}_{H_2O}$ ), which is close to typical ambient water density values in the measurement environment. Also, because one H<sub>2</sub>O value from H<sub>2</sub>O span gas is used, only one parameter in the working equation can be adjusted while others are fixed. Weighing the gain of the working equation more than any other parameter, this parameter is the H<sub>2</sub>O span coefficient (i.e.,  $G_w$ ) in Model (12). The H<sub>2</sub>O span is used to re-estimate  $G_w$  to satisfy the following equation (for more details, see LI-COR Bioscience, 2016):

$$\left| \tilde{\rho}_{H_2O} - \rho_{H_2O}(G_w) \right| = \min \left| \tilde{\rho}_{H_2O} - \rho_{H_2O} \right| \quad (15)$$

After the H<sub>2</sub>O span procedure, the H<sub>2</sub>O gain drift can continue to occur.

Based on the similar considerations as the H<sub>2</sub>O zero drift, the H<sub>2</sub>O gain drift uncertainty is also specified as the maximum range of H<sub>2</sub>O gain drift for the analyzers ( $d_{wg}$ ) that varies with ambient air temperature (Campbell Scientific Inc., 2018b), but normally within the specified range as (see Table 1):

$$d_{wg} = \pm \delta_{H_2O-g} \chi_{H_2O-T} \quad (16)$$

where  $\delta_{H_2O-g}$  is H<sub>2</sub>O gain drift percentage and  $\chi_{H_2O-T}$  is the true H<sub>2</sub>O mixing ratio. This specification is the maximum range of H<sub>2</sub>O measurement uncertainty due to the H<sub>2</sub>O gain drift.

The analyzer performs best, almost without gain drift, when  $T_g$  is equal to the calibration/span ambient air temperature ( $T_c$ , the reason why  $T_c$  still is used here is that zero and span procedures should be performed under similar ambient air temperature conditions). The further  $T_g$  is away from  $T_c$  the more likely it drifts. Using the same way to formulate H<sub>2</sub>O zero drift uncertainty, H<sub>2</sub>O gain drift uncertainty can be approximated for its maximum range as:

$$\Delta \chi_{H_2O}^g = \pm \delta_{H_2O-g} \chi_{H_2O-T} \times \begin{cases} \frac{T_g - T_c}{T_{rh} - T_{rl}} & T_c < T_g < T_{rh} \\ \frac{T_c - T_g}{T_{rh} - T_{rl}} & T_c > T_g > T_{rl} \end{cases} \quad (17)$$

Given the measured value of H<sub>2</sub>O mixing ratio is represented by  $\chi_{H_2O}$ , according to Equation 6, the difference between true and measured H<sub>2</sub>O mixing ratios can be expressed as

$$\chi_{H_2O} - \chi_{H_2O-T} = \Delta \chi_{H_2O}^z + \Delta \chi_{H_2O}^g + \Delta \chi_{H_2O}^s + \Delta \chi_{H_2O}^p \quad (18)$$

From this equation, the true H<sub>2</sub>O mixing ratio is given by:

$$\chi_{H_2O-T} = \chi_{H_2O} - \left( \Delta \chi_{H_2O}^z + \Delta \chi_{H_2O}^g + \Delta \chi_{H_2O}^s + \Delta \chi_{H_2O}^p \right) \quad (19)$$

The term inside the round brackets in this equation is an error term, which generally is smaller, at least, one order than the true value in magnitude. Although the case would not be so true for H<sub>2</sub>O in cold ecosystems (e.g., <math>-5^\circ\text{C}</math>) and/or dry environments, measured value is an appropriate alternative, with the most likelihood, to the true value for the applications of Equation 17. As such,  $\chi_{H_2O-T}$  in Equation 17 can be reasonably approximated by  $\chi_{H_2O}$  for equation applications. Using this approximation, Equation 17 becomes:

$$\Delta \chi_{H_2O}^g = \pm \delta_{H_2O-g} \chi_{H_2O} \times \begin{cases} \frac{T_g - T_c}{T_{rh} - T_{rl}} & T_c < T_g < T_{rh} \\ \frac{T_c - T_g}{T_{rh} - T_{rl}} & T_c > T_g > T_{rl} \end{cases} \quad (20)$$

with  $\chi_{H_2O}$  from measurements, this equation is applicable in estimation for the H<sub>2</sub>O gain drift uncertainty.

### 4.3. $S_c$ and $\Delta\chi_{H_2O}^s$ (Sensitivity-To- $CO_2$ Uncertainty)

Since  $CO_2$  is a weak absorber at the infrared wavelength for  $H_2O$  measurements (i.e., 2.7  $\mu m$ , see Figure 4.7 in Wallace & Hobbs, 2006), it can slightly interfere with  $H_2O$  absorption at this unique wavelength (McDermitt et al., 1993). As such, the power of the same measurement light through several  $H_2O$  gas samples with the same  $H_2O$  molar concentration but different backgrounds of  $CO_2$  amounts would be detected with different values of  $A_w$  for the working equation (see Model [12]). Without the  $S_c$  term in this equation, different  $A_w$  values must result in significantly different  $\rho_{H_2O}$  values although  $\rho_{H_2O}$  is essentially the same. To report the same  $\rho_{H_2O}$  for the air flows with the same  $H_2O$  molar concentration under different  $CO_2$  backgrounds, the different values of  $A_w$  associated with the same  $\rho_{H_2O}$  must be accounted for by  $S_c$  associated with  $A_c$  and  $A_{cs}$  in the working equation (see Model [12]). Similar to  $Z_w$  and  $G_w$  in the equation,  $S_c$  can have an uncertainty in determination of  $\chi_{H_2O}$ . This uncertainty is specified for the gas analyzers by the sensitivity-to- $CO_2$  ( $s_c$ ). For EC155 gas analyzers, the sensitivity-to- $CO_2$  is specified as  $\pm 5 \times 10^{-5}$   $mmolH_2O mol^{-1} (\mu molCO_2 mol^{-1})^{-1}$  (Table 1). As the gas analyzers should be calibrated to produce the minimal uncertainty due to this sensitivity for  $H_2O$  measurements around atmospheric background  $CO_2$  of 415  $\mu molCO_2 mol^{-1}$  (Global Monitoring Laboratory, 2021), the uncertainty in  $\chi_{H_2O}$  measurements due to sensitivity-to- $CO_2$  ( $\Delta\chi_{H_2O}^s$ ) can be quantified as

$$\Delta\chi_{H_2O}^s \equiv s_c (\chi_{CO_2} - 415) \quad 0 \leq \chi_{CO_2} \leq 1000 \mu molCO_2 mol^{-1} \quad (21)$$

Under the atmospheric boundary-layer conditions,  $\chi_{CO_2}$  commonly ranges from 350 to 800  $\mu molCO_2 mol^{-1}$  (LI-Cor Biosciences, 2016). Therefore, using Equation 21,  $\Delta\chi_{H_2O}^s$  satisfies:

$$|\Delta\chi_{H_2O}^s| \leq 585 |s_c| \quad (22)$$

### 4.4. $\Delta\chi_{H_2O}$ ( $H_2O$ Measurement Accuracy)

Substituting Equations 11, 14, 20 and 22 into Equation 10,  $\Delta\chi_{H_2O}$  can be expressed as

$$\Delta\chi_{H_2O} = \pm \left[ 1.96\sigma_{H_2O} + 585|s_c| + (|d_{wz}| + \delta_{H_2O-g}\chi_{H_2O}) \times \begin{cases} \frac{T_g - T_c}{T_{rh} - T_{rl}} & T_c < T_g < T_{rh} \\ \frac{T_c - T_g}{T_{rh} - T_{rl}} & T_c > T_g > T_{rl} \end{cases} \right] \quad (23)$$

This equation is the  $H_2O$  accuracy equation of CPEC systems. It expresses the accuracy of field  $\chi_{H_2O}$  measurements from CPEC infrared gas analyzers in terms of its specifications:  $\sigma_{H_2O}$ ,  $s_c$ ,  $d_{wz}$ , and  $\delta_{H_2O-g}$ ; measured variables  $\chi_{H_2O}$  and  $T_g$ ; and a known variable  $T_c$ . Using this equation and analyzer specification values, the accuracy of field  $H_2O$  measurements can be evaluated as a range.

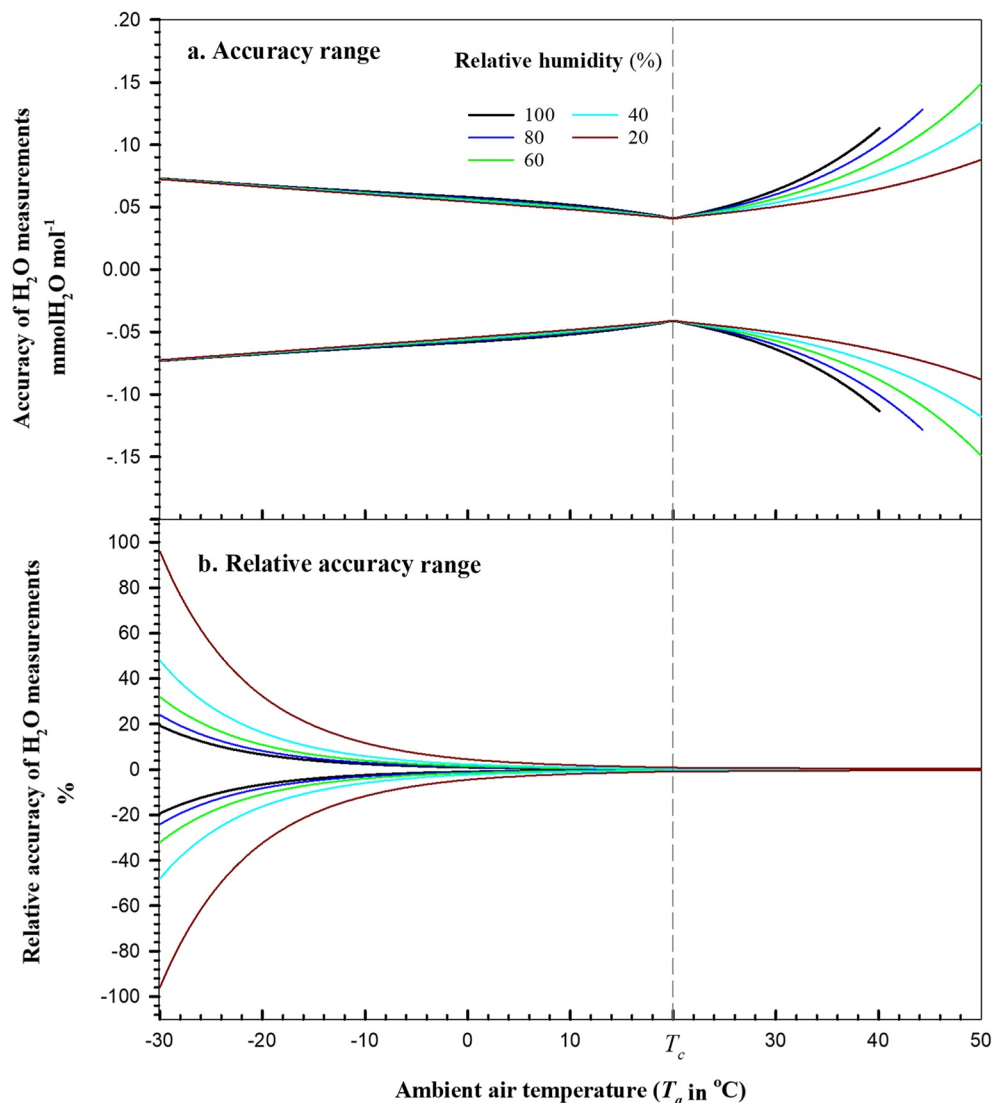
### 4.5. Evaluation on $\Delta\chi_{H_2O}$

Given the values of  $T_g$ ,  $T_c$ , and  $\chi_{H_2O}$  and the analyzer specification values of  $\sigma_{H_2O}$ ,  $s_c$ ,  $d_{wz}$ , and  $\delta_{H_2O-g}$ , the accuracy of  $\chi_{H_2O}$  measurements can be evaluated over a domain of  $T_g$  and  $\chi_{H_2O}$ . To visualize the relationship of accuracy with  $T_g$  and  $\chi_{H_2O}$ , the accuracy is a dependent (the ordinate) of  $T_g$  (the abscissa) at different levels of  $\chi_{H_2O}$ . In addition, this relationship should be evaluated within possible ranges of  $T_g$  and  $\chi_{H_2O}$  that are typically observed in ecosystems.

In practice,  $T_g$  can be approximated with ambient air temperature ( $T_a$ ) in a range of CPEC operations from  $-30$  to  $50^\circ C$ . To evaluate the accuracy under a standard condition instead of a specific field site,  $T_c$  can be set at  $20^\circ C$ , along with atmospheric pressure at 101.325 kPa, as Normal Temperature and Pressure (Wright et al., 2003).

$H_2O$  mixing ratio can be measured by the analyzers from 0 to 79  $mmolH_2O mol^{-1}$ . However, due to the positive dependence of air water vapor saturation on  $T_a$  (Wallace & Hobbs, 2006),  $\chi_{H_2O}$  has a range wider at higher  $T_a$  and narrower at lower  $T_a$ . At  $T_a$  below  $40^\circ C$  under the atmospheric pressure of 101.325 kPa (Wright et al., 2003),  $\chi_{H_2O}$  is lower than 79  $mmolH_2O mol^{-1}$  and, as  $T_a$  decreases, its range becomes narrower and





**Figure 2.** Accuracy of field  $\text{H}_2\text{O}$  measurements from EC155 infrared gas analyzers in closed-path eddy-covariance (CPEC) systems (Campbell Scientific Inc., UT, US) over the operational range of CPEC systems in ambient air temperature from  $-30$  to  $50^{\circ}\text{C}$  under the atmospheric pressure of 101.325 kPa. The vertical dashed line represents ambient air temperature,  $T_c$ , at which the analyzers were calibrated, zeroed, and/or spanned. a. Accuracy of  $\text{H}_2\text{O}$  mixing ratio measurements and b. Relative accuracy of  $\text{H}_2\text{O}$  mixing ratio measurements (i.e., the ratio of accuracy to  $\text{H}_2\text{O}$  mixing ratio).

narrower to be 0 to  $0.38 \text{ mmolH}_2\text{O mol}^{-1}$  at  $-30^{\circ}\text{C}$ . To present the accuracy over the same relative range of air moisture even at different  $T_a$ , the saturation water vapor pressure is used to scale air moisture to 20%, 40%, 60%, 80%, and 100% (i.e., relative humidity [RH]). For each scaled RH value,  $\chi_{\text{H}_2\text{O}}$  can be calculated at different  $T_a$  and  $P$  (Appendix A) for the curves of  $\text{H}_2\text{O}$  accuracy equation with equal RH in Figure 2a.

At  $T_a = T_c$ , the  $\chi_{\text{H}_2\text{O}}$  accuracy is the best at its narrowest range and its magnitude is equal to the sum of precision and sensitivity-to- $\text{CO}_2$  uncertainties ( $<0.0410 \text{ mmolH}_2\text{O mol}^{-1}$  in magnitude). However, away from  $T_c$ , its non-linear range becomes wider, gradually below this  $T_c$  value but more abruptly above, because, as  $T_a$  increases,  $\chi_{\text{H}_2\text{O}}$  at the same RH increases exponentially (Equations A1 and A3 in Appendix A) while  $\Delta\chi_{\text{H}_2\text{O}}$  increases linearly with  $\chi_{\text{H}_2\text{O}}$  in the  $\text{H}_2\text{O}$  accuracy Equation 23. For a case of  $T_c$  at  $20^{\circ}\text{C}$ , the range can be summarized as widest to be  $\pm 0.11 \text{ mmolH}_2\text{O mol}^{-1}$  below  $40^{\circ}\text{C}$  and to be  $\pm 0.15 \text{ mmolH}_2\text{O mol}^{-1}$  above (Figure 2a and  $\text{H}_2\text{O}$  columns in Table 2). In reference to this wider range ( $\pm 0.15 \text{ mmolH}_2\text{O mol}^{-1}$ ), the poorest overall accuracy of  $\text{H}_2\text{O}$  measurements from our study systems can be specified as  $\pm 0.15 \text{ mmolH}_2\text{O mol}^{-1}$ .

**Table 2**

Accuracy of CO<sub>2</sub>/H<sub>2</sub>O Measurements From EC155 Infrared Gas Analyzers in Closed-Path Eddy-Covariance Systems (Campbell Scientific Inc., UT, US) on the Major Values of Background in Ambient Air Temperature, CO<sub>2</sub>, and H<sub>2</sub>O in Ecosystems (Atmospheric Pressure: 101.325 kPa. Calibration Ambient Air Temperature: 20°C)

Ambient air temperature °C	CO <sub>2</sub>				H <sub>2</sub> O			
	415 μmolCO <sub>2</sub> mol <sup>-1</sup>		1,000 μmolCO <sub>2</sub> mol <sup>-1</sup>		60% relative humidity		Saturated	
	Accuracy±	Relative accuracy±	Accuracy±	Relative accuracy±	Accuracy±	Relative accuracy±	Accuracy±	Relative accuracy±
	μmolCO <sub>2</sub> mol <sup>-1</sup>	%	μmolCO <sub>2</sub> mol <sup>-1</sup>	%	mmolH <sub>2</sub> O mol <sup>-1</sup>	%	mmolH <sub>2</sub> O mol <sup>-1</sup>	%
-30	0.7409	0.18	N/A <sup>a</sup>		0.0727	32.12	0.0730	19.34
-25	0.6962	0.17			0.0698	18.52	0.0702	11.18
-22	0.6694	0.16			0.0681	13.44	0.0686	8.12
-20	0.6515	0.16			0.0669	10.89	0.0675	6.59
-18	0.6336	0.15			0.0658	8.85	0.0665	5.36
-15	0.6068	0.15			0.0642	6.52	0.0650	3.96
-10	0.5621	0.14			0.0615	3.97	0.0627	2.43
-6	0.5264	0.13			0.0594	2.70	0.0608	1.66
-5	0.5174	0.12			0.0589	2.46	0.0604	1.51
-2	0.4906	0.12			0.0573	1.85	0.0590	1.14
0	0.4728	0.11			0.0562	1.54	0.0581	0.95
2	0.4549	0.11			0.0551	1.31	0.0570	0.81
5	0.4281	0.10	0.5378	0.05	0.0533	1.02	0.0553	0.63
6	0.4191	0.10	0.5215	0.05	0.0527	0.94	0.0547	0.58
10	0.3834	0.09	0.4565	0.05	0.0500	0.68	0.0519	0.42
15	0.3387	0.08	0.3753	0.04	0.0461	0.45	0.0474	0.28
18	0.3119	0.08	0.3265	0.03	0.0432	0.35	0.0438	0.21
20	0.2940	0.07	0.2940	0.03	0.0410	0.29	0.0410	0.17
22	0.3119	0.08	0.3265	0.03	0.0435	0.27	0.0443	0.16
25	0.3387	0.08	0.3753	0.04	0.0477	0.25	0.0502	0.16
30	0.3834	0.09	0.4565	0.05	0.0569	0.22	0.0637	0.15
35	0.4281	0.10	0.5378	0.05	0.0698	0.20	0.0835	0.14
38	0.4549	0.11	0.5865	0.06	0.0799	0.20	0.0996	0.14
40	0.4728	0.11	0.6190	0.06	0.0879	0.19	0.1126	0.14
42	0.4906	0.12	0.6515	0.07	0.0970	0.19	N/A <sup>b</sup>	
45	0.5174	0.12	0.7003	0.07	0.1133	0.19		
50	0.5621	0.14	0.7815	0.08	0.1489	0.19		

<sup>a</sup>CO<sub>2</sub> mixing ratio is assumed to be lower than 1,000 μmolCO<sub>2</sub> mol<sup>-1</sup> in ambient air temperature below 5°C in ecosystems. <sup>b</sup>H<sub>2</sub>O mixing ratio in saturated air above 40°C under atmospheric pressure of 101.325 kPa is out of the EC155 measurement range (0–79 mmolH<sub>2</sub>O mol<sup>-1</sup>).

Figure 2b shows a pattern of relative accuracy of H<sub>2</sub>O measurements with  $T_a$ . Given a RH above 20%, the relative accuracy diverges to almost 100% as  $T_a$  decreases down to -30°C and converges to 0.35% as  $T_a$  increases up to 50 °C. Given the magnitude of accuracy is in a small order (i.e., <0.15 mmolH<sub>2</sub>O mol<sup>-1</sup>), the divergent pattern is shaped by the exponential decrease in  $\chi_{H_2O}$  saturation amount as  $T_a$  decreases and the convergent pattern is shaped by the exponential increase in  $\chi_{H_2O}$  saturation amount as  $T_a$  increases (see Appendix A). At any  $T_a$  value, the relative accuracy range also can be wide if RH is near zero (i.e., very dry conditions). In ecosystems, unlike CO<sub>2</sub>, H<sub>2</sub>O naturally varies “relatively” large across three orders more in a magnitude

(e.g., from 0.04 mmolH<sub>2</sub>O mol<sup>-1</sup>, when RH is 10% at -30°C under the atmospheric pressure of 101.325 kPa, to 59 mmolH<sub>2</sub>O mol<sup>-1</sup>, when dew point temperature is 35°C under the same atmospheric pressure [National Weather Service, 2021]). Under cold and/or dry conditions, the minimum H<sub>2</sub>O amount could be several orders smaller than the measurement precision of the gas analyzers. In this case, the relative accuracy would be very large and, without specifying the H<sub>2</sub>O status on the measurement background, would not be an appropriate measure to specify the uncertainty of H<sub>2</sub>O measurements from gas analyzers. Accordingly, an unconditional specification of relative accuracy for H<sub>2</sub>O measurements from infrared gas analyzers would mislead users.

To narrow both accuracy and relative accuracy ranges for H<sub>2</sub>O measurements in a lower range of  $T_a$  or  $\chi_{H_2O}$ , frequent zero procedures are needed. Both ranges in Figure 2 are first separately maximized by this study from all uncertainty sources, which is a more solid base for error analysis in H<sub>2</sub>O data applications.

## 5. Accuracy of CO<sub>2</sub> Mixing Ratio Measurements

Accuracy model (9) defines the accuracy of field CO<sub>2</sub> measurements from gas analyzers ( $\Delta\chi_{CO_2}$ ) as

$$\Delta\chi_{CO_2} \equiv \pm \left( \left| \Delta\chi_{CO_2}^z \right| + \left| \Delta\chi_{CO_2}^g \right| + \left| \Delta\chi_{CO_2}^s \right| + \left| \Delta\chi_{CO_2}^p \right| \right) \quad (24)$$

where  $\Delta\chi_{CO_2}^z$  is CO<sub>2</sub> zero drift uncertainty,  $\Delta\chi_{CO_2}^g$  is CO<sub>2</sub> gain drift uncertainty,  $\Delta\chi_{CO_2}^s$  is sensitivity-to-H<sub>2</sub>O uncertainty, and  $\Delta\chi_{CO_2}^p$  is CO<sub>2</sub> precision uncertainty.

### 5.1. $\Delta\chi_{CO_2}^p$ (CO<sub>2</sub> Precision Uncertainty)

Using the same approach for  $\Delta\chi_{H_2O}^p$ ,  $\Delta\chi_{CO_2}^p$  is formulated as:

$$\Delta\chi_{CO_2}^p = \pm 1.96 \times \sigma_{CO_2} \quad (25)$$

### 5.2. $\Delta\chi_{CO_2}^z$ (CO<sub>2</sub> Zero Drift Uncertainty) and $\Delta\chi_{CO_2}^g$ (CO<sub>2</sub> Gain Drift Uncertainty)

The working model of gas analyzers for  $\chi_{CO_2}$  is similar to Model (12) for  $\chi_{H_2O}$  in formulation, according to LI-Cor Biosciences (2016), given by:

$$\rho_{CO_2} = P \sum_{i=1}^5 a_{ci} \left\{ 1 - \left[ \frac{A_c}{A_{cs}} + S_w \left( 1 - \frac{A_w}{A_{ws}} \right) \right] Z_c \right\}^i \left\{ \frac{G_c}{P} \right\}^i \quad (26)$$

$$\chi_{CO_2} = \rho_{CO_2} \left[ \frac{P}{R^* (T_g + 273.15)} - \frac{\rho_{H_2O}}{1000} \right]^{-1}$$

where  $\rho_{CO_2}$  is CO<sub>2</sub> molar concentration ( $\mu\text{molCO}_2 \text{ m}^{-3}$ );  $a_{ci}$  ( $i = 1, 2, 3, 4, \text{ or } 5$ ) is a coefficient of the fifth order polynomial in the terms inside curly brackets;  $S_w$  is the cross-sensitivity of the detector to H<sub>2</sub>O, while detecting CO<sub>2</sub>, in the wavelength for CO<sub>2</sub> measurements (hereinafter referred as sensitivity-to-H<sub>2</sub>O);  $Z_c$  is CO<sub>2</sub> zero adjustment (i.e., CO<sub>2</sub> zero coefficient);  $G_c$  is CO<sub>2</sub> gain adjustment (i.e., CO<sub>2</sub> gain coefficient commonly as CO<sub>2</sub> span coefficient). The parameters of  $a_{ci}$ ,  $Z_c$ ,  $G_c$ , and  $S_w$  in this model are statistically estimated to establish a CO<sub>2</sub> working equation in the production calibration against a series of standard CO<sub>2</sub> gases over the ranges of  $\rho_{H_2O}$  and  $P$  (hereinafter referred as calibration). The CO<sub>2</sub> working equation (i.e., Model [26] with estimated parameters) is used inside the gas analyzer to compute  $\rho_{CO_2}$  and  $\chi_{CO_2}$  from field measurements of  $A_c$ ,  $A_{cs}$ ,  $A_w$ ,  $A_{ws}$ ,  $P$ , and  $T_g$ .

Because of similarity in model principals and parameter implications between Models (12) and (26), using the same analyses and rationales as for  $\Delta\chi_{H_2O}^z$  and  $\Delta\chi_{H_2O}^g$ ,  $\Delta\chi_{CO_2}^z$  is formulated as:

$$\Delta\chi_{CO_2}^z \equiv d_{cz} \times \begin{cases} \frac{T_g - T_c}{T_{rh} - T_{rl}} & T_c < T_g < T_{rh} \\ \frac{T_c - T_g}{T_{rh} - T_{rl}} & T_c > T_g > T_{rl} \end{cases} \quad (27)$$

where  $d_{cz}$  is the maximum CO<sub>2</sub> zero drift specified for gas analyzers, and  $\Delta\chi_{CO_2}^s$  is formulated as:

$$\Delta\chi_{CO_2}^s \equiv \pm\delta_{CO_2-g}\chi_{CO_2} \times \begin{cases} \frac{T_g - T_c}{T_{rh} - T_{rl}} & T_c < T_g < T_{rh} \\ \frac{T_c - T_g}{T_{rh} - T_{rl}} & T_c > T_g > T_{rl} \end{cases} \quad (28)$$

where  $\delta_{CO_2-g}$  is the maximum CO<sub>2</sub> gain drift percentage specified for gas analyzers. Both  $d_{cz}$  and  $\delta_{CO_2-g}$  are given in analyzer specifications (Table 1).

### 5.3. $\Delta\chi_{CO_2}^s$ (Sensitivity-To-H<sub>2</sub>O Uncertainty)

Since H<sub>2</sub>O also is a weak absorber at the infrared wavelength for CO<sub>2</sub> measurements (i.e., 4.3  $\mu$ m, Campbell Scientific Inc., 2018a; LI-COR Biosciences 2016), it interferes with CO<sub>2</sub> absorption slightly (McDermitt et al., 1993). As such, the power of identical measurement light through several gas samples with the same CO<sub>2</sub> molar concentration but different backgrounds of H<sub>2</sub>O molar concentrations would result in different values of  $A_c$  into the CO<sub>2</sub> working equation. Without the  $S_w$  term in this equation, different  $A_c$  values will result in significantly different  $\rho_{CO_2}$  values although  $\rho_{CO_2}$  is actually the same. To report the same  $\rho_{CO_2}$  for the air flows with the same CO<sub>2</sub> molar concentration under different H<sub>2</sub>O backgrounds, the different values of  $A_c$  associated with  $\rho_{CO_2}$  must be accounted for by  $S_w$  associated with  $A_w$  and  $A_{ws}$  as shown in Model (26). However,  $S_w$  can have an uncertainty in its accountability. This uncertainty is specified by the sensitivity-to-H<sub>2</sub>O ( $s_w$ ). For EC155 gas analyzers, the sensitivity-to-H<sub>2</sub>O is specified as  $\pm 5.6 \times 10^{-8} \mu\text{molCO}_2 \text{ mol}^{-1} (\text{mmolH}_2\text{O mol}^{-1})^{-1}$  (Table 1). Given that the gas analyzers for CO<sub>2</sub> works best for dry air, the uncertainty in  $\chi_{CO_2}$  measurements due to sensitivity-to-H<sub>2</sub>O ( $\Delta\chi_{CO_2}^s$ ) can be quantified as

$$\Delta\chi_{CO_2}^s \equiv s_w\chi_{H_2O} \quad 0 \leq \chi_{H_2O} \leq 79 \text{ mmolH}_2\text{O mol}^{-1} \quad (29)$$

Accordingly,  $\Delta\chi_{CO_2}^s$  can be reasonably expressed as:

$$|\Delta\chi_{CO_2}^s| \leq 79|s_w| \quad (30)$$

### 5.4. $\Delta\chi_{CO_2}$ (CO<sub>2</sub> Measurement Accuracy)

Substituting Equations 25, 27, 28 and 30 into Equation 24,  $\Delta\chi_{CO_2}$  can be expressed as:

$$\Delta\chi_{CO_2} = \pm \left[ 1.96\sigma_{CO_2} + 79|s_w| + (|d_{cz}| + \delta_{CO_2-g}\chi_{CO_2}) \times \begin{cases} \frac{T_g - T_c}{T_{rh} - T_{rl}} & T_c < T_g < T_{rh} \\ \frac{T_c - T_g}{T_{rh} - T_{rl}} & T_c > T_g > T_{rl} \end{cases} \right] \quad (31)$$

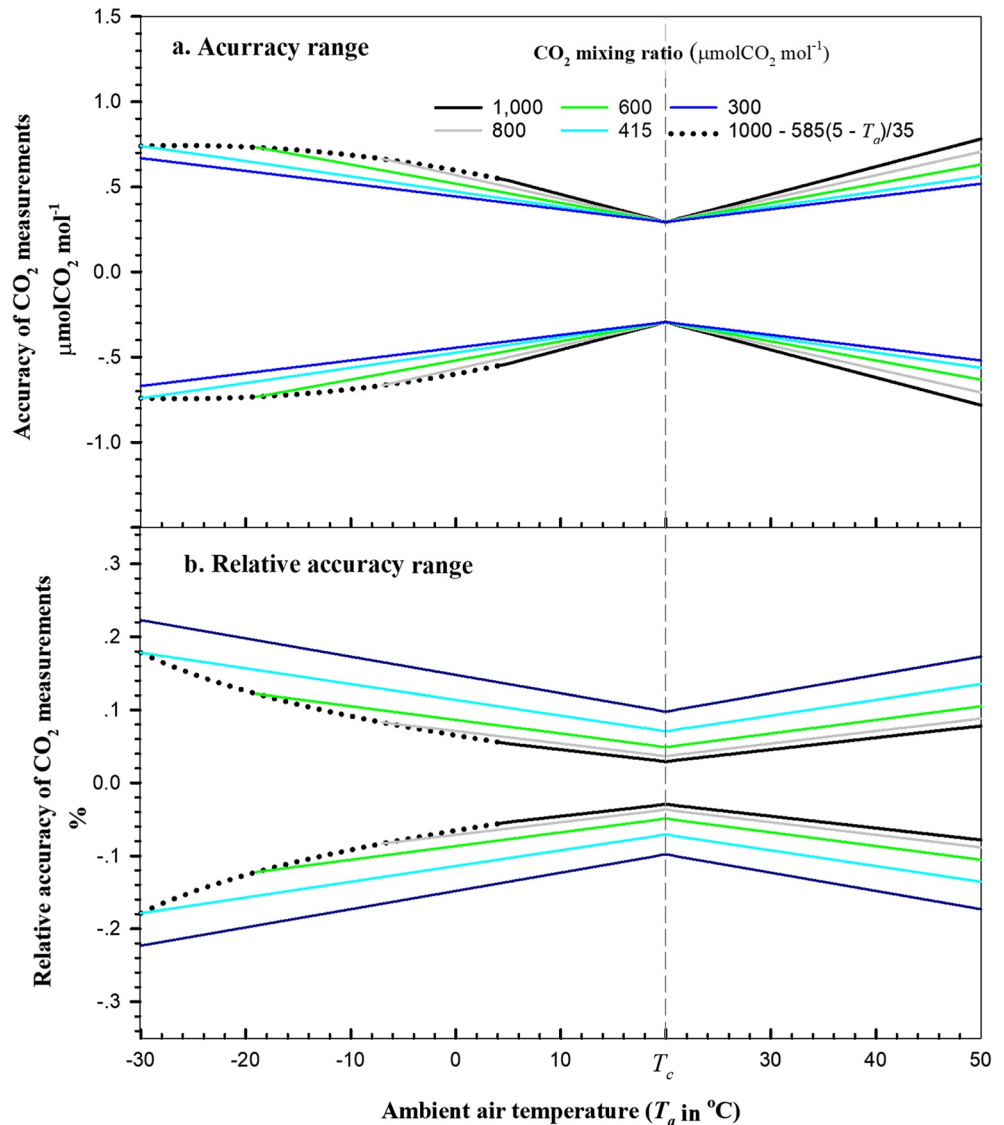
This equation is the CO<sub>2</sub> accuracy equation of infrared gas analyzers in CPEC systems. It expresses the accuracy of  $\chi_{CO_2}$  measurements from the analyzers in terms of their specifications:  $\sigma_{CO_2}$ ,  $s_w$ ,  $d_{cz}$ , and  $\delta_{CO_2-g}$ ; measured variables  $\chi_{CO_2}$  and  $T_g$ ; and a known variable  $T_c$ . Using this equation and analyzer specification values, the accuracy of field CO<sub>2</sub> measurements can be evaluated as a range.

### 5.5. Evaluation on $\Delta\chi_{CO_2}$

Using the CO<sub>2</sub> accuracy Equation 31 along with the analyzer specifications of  $\sigma_{CO_2}$ ,  $s_w$ ,  $d_{cz}$ , and  $\delta_{CO_2-g}$ ,  $\Delta\chi_{CO_2}$  can be evaluated over a domain of  $T_g$  and  $\chi_{CO_2}$  for a given  $T_c$ . Similar to the presentation of H<sub>2</sub>O accuracy in Figure 2,  $\Delta\chi_{CO_2}$  is a dependent (the ordinate) of  $T_g$  (the abscissa) from  $-30$  to  $50^\circ\text{C}$  at different levels of  $\chi_{CO_2}$  within the range typically observed in ecosystems.  $T_c$  is set for the normal temperature of  $20^\circ\text{C}$  (Wright et al., 2003) and  $T_g$  is surrogated by  $T_a$  (Figure 3).

#### 5.5.1. $\chi_{CO_2}$ Range

$\chi_{CO_2}$  is measured by the EC155 infrared CO<sub>2</sub>/H<sub>2</sub>O analyzers up to  $1,000 \mu\text{molCO}_2 \text{ mol}^{-1}$ . In the atmosphere,  $\chi_{CO_2}$  average is currently  $\sim 415 \mu\text{molCO}_2 \text{ mol}^{-1}$  (Global Monitoring Laboratory, 2021). However, in



**Figure 3.** Accuracy of field CO<sub>2</sub> measurements from the EC155 infrared gas analyzers in closed-path eddy-covariance systems over the operational ranges of ambient air temperature (−30° to 50°C) at 101.325 kPa atmospheric pressure. The vertical dashed line represents ambient air temperature, T<sub>c</sub>, at which an analyzer was calibrated, zeroed, and/or spanned. When T<sub>a</sub> decreases from 5° to −30°C, CO<sub>2</sub> mixing ratio on black dotted lines linearly decreases from 1,000 μmolCO<sub>2</sub> mol<sup>-1</sup> at 5°C to 415 μmolCO<sub>2</sub> mol<sup>-1</sup> (i.e., atmospheric background value) at −30°C. (a) Accuracy of CO<sub>2</sub> mixing ratio measurements and (b) Relative accuracy of CO<sub>2</sub> mixing ratio measurements (i.e., the ratio of accuracy to CO<sub>2</sub> mixing ratio).

terrestrial ecosystems, where the analyzers are deployed,  $\chi_{CO_2}$  fluctuates with human induced emissions and biological processes such as plant physiological metabolism, soil microorganism respiration, and animal physiological activities (Wang et al., 2016); aerodynamic regimes such as wind speed, wind direction related to terrain topography (de Araújo et al., 2010), and vertical wind gradient (Yang et al., 2007); and thermodynamic states such as air temperature, soil temperature, and boundary-layer stability (Ohkubo et al., 2008). As discussed above,  $\chi_{CO_2}$  in ecosystems commonly ranges from 350 to 800 μmolCO<sub>2</sub> mol<sup>-1</sup> depending on biological processes, aerodynamic regimes, and thermodynamic states. This range is extended from 300 to 1,000 μmolCO<sub>2</sub> mol<sup>-1</sup> as the possible range within which  $\Delta\chi_{CO_2}$  is evaluated. Because of the dependence of  $\Delta\chi_{CO_2}$  on  $\chi_{CO_2}$  (Equations 28 and 31), to show the accuracy at different levels of  $\chi_{CO_2}$ , the range is further divided into five grades of 300, 415 (atmospheric background value), 600, 800, and 1000 μmolCO<sub>2</sub> mol<sup>-1</sup> for evaluation presentations as in Figure 3a.

However, in terrestrial ecosystems, with decreasing  $T_a$  from its plant physiological threshold for growth and development, biological processes diminish. Generally, while  $T_a$  decreases from 5°C to negative, plants and microorganisms gradually go dormant or finish their life spans if frozen (Taiz et al., 2014). In such a process, however, soil temperature related to microorganism respiration and/or activities in deeper layers decreases in lag (Widén & Majdi, 2001). While  $T_a$  is below 0°C, soil temperature at some depths is likely still above freezing, and soil microorganisms are still active (Rosenberg et al., 1983). Due to the activities sometimes under the conditions of low wind along with atmospheric stable stratification,  $\chi_{CO_2}$  in ecosystems may still be higher than the  $CO_2$  background concentration (Nicolini et al., 2018), but it should approach this background concentration at very low  $T_a$ . The value of  $-30^\circ C$  at the low end of the specified temperature range for CPEC operations can be considered lower than enough. While  $T_a$  decreases to  $-30^\circ C$ , the  $\chi_{CO_2}$  in ecosystems, if higher, should gradually decrease to the  $CO_2$  background value. Accordingly,  $\chi_{CO_2}$ , if higher, should start at 5°C to converge asymptotically to  $415 \mu mol CO_2 mol^{-1}$  at  $-30^\circ C$ . Without the asymptotical function for the convergence boundary of curve trend, the convergence can be conservatively assumed as a simple linear trend from 1,000 (i.e., maximum) to  $415 \mu mol CO_2 mol^{-1}$  as the boundary while  $T_a$  decreases from 5° to  $-30^\circ C$ . The  $CO_2$  measurement accuracy,  $\Delta\chi_{CO_2}$ , at each  $CO_2$  grade is evaluated up to the boundary on the dotted trend curve as shown in Figure 3.

### 5.5.2. $\Delta\chi_{CO_2}$ Evaluation

At  $T_a = T_c$ , the  $\chi_{CO_2}$  accuracy is the best at its narrowest range and its magnitude is the sum of precision and sensitivity-to- $H_2O$  uncertainties ( $0.29 \mu mol CO_2 mol^{-1}$ ). However, away from  $T_c$ , its range expands in a near-linear fashion. In case of  $T_c$  at  $20^\circ C$ , the  $\Delta\chi_{CO_2}$  range can be summarized as a maximum to be  $\pm 0.78 \mu mol CO_2 mol^{-1}$  in ecosystems at the extreme conditions (e.g.,  $50^\circ C$ . Figure 3a and  $CO_2$  columns in Table 2). The relative  $CO_2$  accuracy has its maximum range of  $\pm 0.23\%$  (Figure 3b).

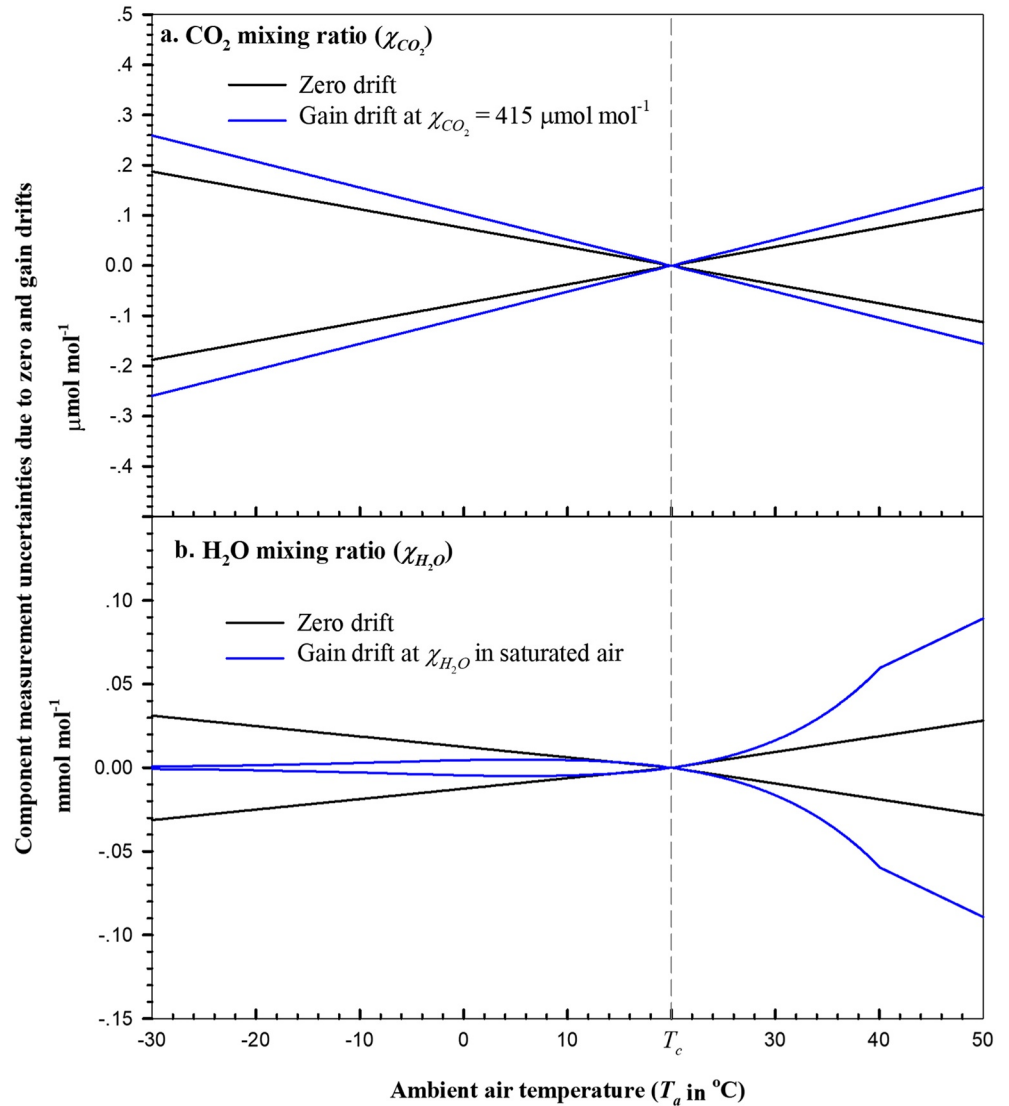
Same as Figure 2, both ranges in Figures 3a and 3b are first separately maximized by this study from all uncertainty sources, which is a more solid base for error analyses in  $CO_2$  data applications.

## 6. Discussion

The primary objective of this study is to quantify the accuracy of field  $CO_2/H_2O$  measurements by infrared gas analyzers in CPEC systems from their specifications for their zero drift, gain drift, sensitivity-to- $CO_2/H_2O$ , and precision. In practice, the accuracy is used to quantify the uncertainty of  $CO_2/H_2O$  data for applications. Additionally, the relationship of accuracy to different specification terms (e.g., gain drift) as component uncertainties can be considered as rationale to guide the analyzer field maintenance in maximally improving field measurement accuracy.

### 6.1. Application Example

As discussed in the introduction,  $\chi_{H_2O}$  along with  $T_s$  is applicable to calculate  $T_a$  (Kaimal & Gaynor, 1991). Given the function,  $T_a(\chi_{H_2O}, T_s)$ , from a theoretical basis of first principles, the equation itself does not have any error and the calculated  $T_a$  should be accurate as long as the values of  $\chi_{H_2O}$  and  $T_s$  are exact. For our subject, however,  $\chi_{H_2O}$  and  $T_s$  are measured by the CPEC systems deployed in the field under changing environments. Their measured values must include measurement uncertainty in  $\chi_{H_2O}$ , denoted by  $\Delta\chi_{H_2O}$  (i.e., field  $\chi_{H_2O}$  measurement accuracy), and in  $T_s$  as well, denoted by  $\Delta T_s$  (i.e., field  $T_s$  measurement accuracy).  $\Delta\chi_{H_2O}$  and/or  $\Delta T_s$  unavoidably propagate to the calculated  $T_a$  through equation output as uncertainty, denoted by  $\Delta T_a$ . In numerical analysis (Burden & Faires, 1993) or in statistics (Snedecor & Cochran, 1989), any applicable equation requires the specification for its uncertainty term. Therefore,  $T_a(\chi_{H_2O}, T_s)$  should include a specification of the respective accuracy expressed as uncertainty bounds, which are the maximum and minimum limits of calculated  $T_a$  for any given pair of  $\chi_{H_2O}$  and  $T_s$ . This accuracy is needed by any application. It should be specified through the relationship of  $\Delta T_a$  to  $\Delta\chi_{H_2O}$  and  $\Delta T_s$ .



**Figure 4.** Component measurement uncertainties due to the zero and gain drifts of EC155 infrared gas analyzers in closed-path eddy-covariance systems over the operational range of ambient air temperature under the atmospheric pressure of 101.325 kPa. (a) Zero and gain drift uncertainties in  $\chi_{CO_2}$  measurements and (b) Zero and gain drift uncertainties in  $\chi_{H_2O}$  measurements. Dashed line represents ambient air temperature,  $T_c$ , at which an analyzer was calibrated, zeroed, and/or spanned.

As field measurement accuracies, both  $\Delta\chi_{H_2O}$  and  $\Delta T_s$  can be reasonably considered as small increments in a calculus sense. Following the principles of differential equation and considering that  $\chi_{H_2O}$  and  $T_s$  are measured from two independent sensors,  $\Delta T_a$  is a total differential of  $T_a(\chi_{H_2O}, T_s)$  with respect to  $\chi_{H_2O}$  and  $T_s$ , given by:

$$\Delta T_a = \frac{\partial T_a(\chi_{H_2O}, T_s)}{\partial \chi_{H_2O}} \Delta \chi_{H_2O} + \frac{\partial T_a(\chi_{H_2O}, T_s)}{\partial T_s} \Delta T_s \quad (32)$$

Apparently,  $\Delta\chi_{H_2O}$  is a required term for evaluation of  $\Delta T_a$ . In this equation, the two partial derivatives can be derived from the explicit function of  $T_a(\chi_{H_2O}, T_s)$ ,  $\Delta\chi_{H_2O}$  is estimated by the application of this study, and  $\Delta T_s$  can be acquired from the specification for 3D sonic anemometers (Zhou et al., 2018). With the two partial derivatives,  $\Delta\chi_{H_2O}$ , and  $\Delta T_s$ ,  $\Delta T_a$  can be evaluated as a function of  $\chi_{H_2O}$  and  $T_s$ .

## 6.2. Rationale to Guide Analyzer Field Maintenance

An infrared gas analyzer works at its best in the same environment as its manufacturing conditions, which is shown in Figures 2 and 3 for measurement accuracies associated with  $T_c$ . The analyzer works better while  $T_a$  is closer to  $T_c$ . However, in practice, it is deployed commonly for long-term measurements in the field under changing weather conditions through seasonal climates. Most of the time these conditions are different from those in the manufacturing process. Over time, the analyzer gradually drifts to some degree, although the drifts should be within its specifications, and needs maintenance.

As discussed in Sections 4 and 5, the uncertainty in analyzer measurements from the sources of precision and sensitivity-to- $\text{CO}_2/\text{H}_2\text{O}$  uncertainties is not improvable through field maintenance, but it is small ( $\pm 0.29 \mu\text{molCO}_2 \text{ mol}^{-1}$  for  $\text{CO}_2$ , Equations 25, 30 and 31, and  $\pm 0.04 \text{ mmolH}_2\text{O mol}^{-1}$  for  $\text{H}_2\text{O}$ , Equations 11, 22 and 23). However, the zero and gain drifts of analyzers are major components of uncertainty in determination of their field measurement accuracies (Figures 2–4, Equations 14, 20, 27 and 28). Fortunately, the drifts are adjustable through the zero and span procedures (see Sections 4.1 and 4.2). The procedures are recommended to minimize the drift influence on field measurement uncertainties. Therefore, the manufacturers of infrared gas analyzers designed software and hardware tools for the procedures and recommended the procedure schedule using the tools (Campbell Scientific Inc., 2018a; LI-COR Biosciences, 2016). As shown in Figures 2–4, this study helps users understand, through visualization, how to schedule, perform, and assess the procedures.

### 6.2.1. $\text{CO}_2$ Zero and Span Procedures

As shown in Figures 3 and 4a, the  $\text{CO}_2$  zero drift and/or  $\text{CO}_2$  gain drift can bring appreciable uncertainties to field  $\text{CO}_2$  measurements in a similar trend over the full  $T_a$  range within which CPEC systems operate. Both drifts should be adjusted near  $T_a$  around which a corresponding analyzer runs. Fortunately, unlike the  $\text{H}_2\text{O}$  gas,  $\text{CO}_2$  gas can be conveniently used under any  $T_a$  environment. Referring to Figure 4a, the zero and gain drifts for field  $\text{CO}_2$  measurements should be adjusted, through the zero and span procedures, at  $T_c$  close to the daily mean of  $T_a$ . A CPEC310 system automatically zeroes and spans its analyzer for field  $\text{CO}_2$  measurements in a timing interval set by a user at a resolution of a minute (Campbell Scientific Inc., 2018b). According to the range of  $T_a$  daily cycle, the procedures are set around its daily mean. Given that  $T_a$  fluctuates within a daily range much narrower than  $40^\circ\text{C}$ , the CPEC system can run at  $T_a$  within  $\pm 20$  of  $T_c$ . In this way after the procedures, for the case in atmospheric background of  $\text{CO}_2$  (i.e.,  $415 \mu\text{molCO}_2 \text{ mol}^{-1}$ ), the widest possible range of  $\pm 0.74 \mu\text{molCO}_2 \text{ mol}^{-1}$  (see the left  $\text{CO}_2$  accuracy column in Table 2) for field  $\text{CO}_2$  measurement accuracy can be narrowed by 36% to  $\pm 0.47 \mu\text{molCO}_2 \text{ mol}^{-1}$  (estimated according to the  $\text{CO}_2$  accuracy at  $T_a$  of  $20^\circ\text{C}$  above and below calibration ambient air temperature), which would be significant improvement in  $\text{CO}_2$  measurements from CPEC automatic  $\text{CO}_2$  zero and span procedures.

Additionally, in ecosystems during the growing season, the gain drift likely causes more uncertainty than zero drift due to more events of higher  $\text{CO}_2$  mixing ratio from active respiration of soil-plant continuum under stable atmospheric stratification (Widén & Majdi., 2001; Yang et al., 2007). In such a season, more frequent  $\text{CO}_2$  span procedures are preferred as if a CPEC310 system (Campbell Scientific Inc., UT, US) ran with automatic zero and  $\text{CO}_2$  span procedures at a timing interval optionally set by a user.

### 6.2.2. $\text{H}_2\text{O}$ Zero and Span Procedures

As shown in Figure 4b, the uncertainty in  $\text{H}_2\text{O}$  measurements from the zero drift increases with  $T_a$  away from  $T_c$  in the same way as for  $\text{CO}_2$  measurements. Therefore,  $\text{H}_2\text{O}$  zero procedure should be performed in the same way as for  $\text{CO}_2$  zero procedure. However, the uncertainty from the gain drift exponentially diverges, as  $T_a$  increases away from  $T_c$ , up to  $\pm 0.0892 \text{ mmolH}_2\text{O mol}^{-1}$  at  $50^\circ\text{C}$  (from data for Figure 4b) and gradually converges, as  $T_a$  decreases away from  $T_c$ , down to a small magnitude,  $\pm 0.0007 \text{ mmolH}_2\text{O mol}^{-1}$ , at  $-30^\circ\text{C}$  (from data for Figure 4b). The exponential divergence above  $T_c$  is formed by the exponential increase in air  $\text{H}_2\text{O}$  saturation with  $T_a$  increase (Buck, 1981) while the  $\text{H}_2\text{O}$  gain drift uncertainty linearly increases with  $\lambda_{\text{H}_2\text{O}}$  as  $T_a$  increases away from  $T_c$  (Equation 20). Below  $T_c$ , the gradual convergence is formed by the exponential decrease in air  $\text{H}_2\text{O}$  saturation with  $T_a$  decrease although the  $\text{H}_2\text{O}$  gain drift uncertainty still linearly increases with  $\lambda_{\text{H}_2\text{O}}$  as  $T_a$  decreases away from  $T_c$  (Equation 20). This trend of  $\text{H}_2\text{O}$  gain drift uncer-



tainty with  $T_a$  is a rationale to judge the need of the H<sub>2</sub>O span procedure under different environments to adjust the H<sub>2</sub>O gain drift.

Given  $T_c$  at 20°C (i.e., the H<sub>2</sub>O span procedure is performed at 20°C), the uncertainty from gain drift for saturated air at  $T_a$  below 0°C is estimated within  $\pm 0.0046$  mmolH<sub>2</sub>O mol<sup>-1</sup> (from data for Figure 4b). Given  $T_c$  equal to 5°C, this estimated uncertainty would be within  $\pm 0.0012$  mmolH<sub>2</sub>O mol<sup>-1</sup> equal to 5% H<sub>2</sub>O zero drift uncertainty,  $\pm 0.0219$  mmolH<sub>2</sub>O mol<sup>-1</sup>, if not zeroed at  $T_a$  below 5°C. In the range of  $T_a$  below 0°C, considering the small H<sub>2</sub>O gain drift uncertainty and inconvenient application of H<sub>2</sub>O span gas from any dew point generator (LI-COR Biosciences, 2004), the H<sub>2</sub>O span procedure may be unnecessary, but the H<sub>2</sub>O zero procedure becomes critically necessary in narrowing H<sub>2</sub>O measurement accuracy. This recommendation eases the users who are worried about H<sub>2</sub>O measurement accuracy in the cold seasons during which the H<sub>2</sub>O span procedure for analyzers is hardly performed in a convenient way while the H<sub>2</sub>O zero procedure still can be automatically performed (Campbell Scientific Inc. 2018a). In contrast, in the higher  $T_a$  range (e.g., 30°C, see Figure 2a and Table 2) under humid conditions, this H<sub>2</sub>O span procedure must be more frequently needed. Unfortunately, unlike the CO<sub>2</sub> gas, H<sub>2</sub>O gas cannot be conveniently used for automatic H<sub>2</sub>O span procedure under any  $T_a$  environment (LI-COR Biosciences, 2004). Nonetheless, the results from this study emphasize the need for the H<sub>2</sub>O span procedure for the gas analyzers at a higher  $T_a$  range in humid conditions.

Given that the H<sub>2</sub>O and CO<sub>2</sub> zero drifts are automatically adjusted together through all-zero procedure and the H<sub>2</sub>O gain drift is manually adjusted through the H<sub>2</sub>O span procedure in a manual-recommended interval (Campbell Scientific Inc., 2018a; LI-COR Biosciences, 2016) practically over a range of  $T_a$  from 5° to 30°C, the widest range,  $\pm 0.15$  mmolH<sub>2</sub>O mol<sup>-1</sup>, of H<sub>2</sub>O measurement accuracy in case of  $T_c$  at 20°C (see the H<sub>2</sub>O accuracy column for 60% relative humidity in Table 2) can be narrowed by 27% to  $\pm 0.11$  mmolH<sub>2</sub>O mol<sup>-1</sup> (estimated by setting  $T_c$  as 30°C). While warranting H<sub>2</sub>O measurement accuracy, CPEC automatic zero procedure over the system operational range of  $T_a$  and the manual H<sub>2</sub>O span over a practical range of  $T_a$  also can improve the accuracy as analyzed from Figure 2 and Table 2.

### 6.3. Benefit From the Automatic Zero Procedure

Looking at the saturated H<sub>2</sub>O columns in Table 2, the poorest accuracy of H<sub>2</sub>O measurements at  $T_a$  below 0°C is  $\pm 0.073$  mmolH<sub>2</sub>O mol<sup>-1</sup>. In the same  $T_a$  range, automatically adjusting the zero drift on a daily base for the gas analyzer to run at  $T_a$  within  $\pm 20^\circ\text{C}$  of  $T_c$ , the zero procedure narrows this range by 36% at least to  $\pm 0.047$  mmolH<sub>2</sub>O mol<sup>-1</sup> (estimated by setting  $T_c$  to be  $-10^\circ\text{C}$ ). At the same time, the range of relative accuracy (Figure 2b) can be narrowed by the same percentage. Apparently, while the H<sub>2</sub>O span procedure is unnecessary, the accuracy and relative accuracy of H<sub>2</sub>O measurements at  $T_a$  below 0°C can be greatly benefited from the automatic zero procedure.

## 7. Concluding Remarks

As advanced by the International Organization for Standardization (2012), measurement accuracy is defined as a combination of trueness (systematic uncertainty) and precision (random uncertainty), being a range of the composite uncertainty from all measurement uncertainty sources. Adopting this definition, the accuracy of field CO<sub>2</sub>/H<sub>2</sub>O measurements from the infrared gas analyzers in CPEC systems is defined as a maximum range of composite measurement uncertainty sourced from trueness (zero drift, gain drift, and sensitivity-to-CO<sub>2</sub>/H<sub>2</sub>O uncertainties) and precision (random measurement uncertainty). The analyzers are specified for their measurement performances in terms of zero drift, gain drift, sensitivity-to-CO<sub>2</sub>/H<sub>2</sub>O, and precision (Campbell Scientific Inc., 2018a). The uncertainty terms are comprehended into accuracy model (10) as a composite uncertainty to describe the accuracy of field CO<sub>2</sub>/H<sub>2</sub>O measurements from the analyzers. Based on instrumentation technology, the specified uncertainty terms of analyzers calibrated under manufacturer ambient conditions are incorporated into the model. The incorporation formulates CO<sub>2</sub> accuracy Equation 31 and H<sub>2</sub>O accuracy Equation 23. According to atmospheric physics and environmental parameters, the equations are used to evaluate the analyzers for their accuracies of field CO<sub>2</sub> and H<sub>2</sub>O measurements in ecosystems over the analyzer operational range of ambient air temperature ( $T_a$ ) under the normal atmospheric pressure of 101.325 kPa (Wright et al., 2003).

The maximum accuracy range of field CO<sub>2</sub> measurements at 415 μmolCO<sub>2</sub> mol<sup>-1</sup> (atmospheric background) from the analyzers is ±0.74 μmolCO<sub>2</sub> mol<sup>-1</sup> (relatively ±0.18%) and at 1,000 μmolCO<sub>2</sub> mol<sup>-1</sup>, assumed as the highest CO<sub>2</sub> in ecosystems, is ±0.78 μmolCO<sub>2</sub> mol<sup>-1</sup> (relatively ±0.08%, Figure 3 and Table 2). Accordingly, for the CO<sub>2</sub> measurements in ecosystems from our example gas analyzers, the overall accuracy can be specified as ±0.78 μmolCO<sub>2</sub> mol<sup>-1</sup>, relatively being within ±0.18%.

The maximum accuracy range of field H<sub>2</sub>O measurements from the gas analyzers can be specified as ±0.15 mmol-H<sub>2</sub>O mol<sup>-1</sup> (Figure 2a and Table 2). For the H<sub>2</sub>O measurements in ecosystems, relative accuracy is not recommended to specify the performance of analyzers, because, without knowing the air moisture level in magnitude for measurements, the relative accuracy in percent for H<sub>2</sub>O measurements is not practically descriptive.

Accuracy Equations 23 and 31 are not only applicable for error/uncertainty analyses in CO<sub>2</sub>/H<sub>2</sub>O data applications (see Section 6.1), but also used as rationales to guide and assess the field maintenances on the gas analyzers. Equation 31, as shown in Figures 3a and 4a, guides users to adjust the CO<sub>2</sub> zero drift through the all-zero (i.e., CO<sub>2</sub>/H<sub>2</sub>O zero) procedure and the CO<sub>2</sub> gain drift through the CO<sub>2</sub> span procedure. Given that  $T_a$  fluctuates within a daily range much narrower than 40 °C, the automatic zero/span procedures could ensure that infrared gas analyzers run at  $T_a$  reasonably within ±20 of  $T_c$  at which an analyzer is calibrated or zeroed/spanned. As assessed, on the atmospheric CO<sub>2</sub> background, the procedures can narrow the maximum CO<sub>2</sub> accuracy range by 36% from ±0.74 to ±0.47 μmolCO<sub>2</sub> mol<sup>-1</sup>.

Equation 23 as shown in Figures 2a and 4b guides users to adjust the H<sub>2</sub>O zero drift in the same way as for CO<sub>2</sub>. It rationalizes that the H<sub>2</sub>O gain drift needs to be adjusted with more attention to the higher  $T_a$  range (e.g., above 30°C) under humid conditions. This adjustment is not necessary below 0°C. Given that the H<sub>2</sub>O zero drift is automatically adjusted daily through the all-zero procedure, the H<sub>2</sub>O gain drift can be manually adjusted through the H<sub>2</sub>O span procedure in a manual-recommended interval while  $T_a$  is in an operational range of 5°–30°C for H<sub>2</sub>O span (i.e.,  $T_c$  between 5° and 30°C). Such a zero/span protocol could narrow the H<sub>2</sub>O accuracy range of ±0.15 by 27% to ±0.11 mmolH<sub>2</sub>O mol<sup>-1</sup>.

In the  $T_a$  range below 0°C, the zero procedure can narrow the H<sub>2</sub>O accuracy range of ±0.073 mmol-H<sub>2</sub>O mol<sup>-1</sup>, in case of air saturation, by 36% at least to ±0.047 mmolH<sub>2</sub>O mol<sup>-1</sup>. At the same time, the range of relative accuracy (Figure 2b) can be narrowed by the same percentage. For cold environments, the unnecessary for H<sub>2</sub>O span procedures and greater necessity for automatic all-zero procedures are first addressed in this study. The former relieves users from worrying about the H<sub>2</sub>O measurement uncertainty in freezing conditions from the H<sub>2</sub>O gain drift, which is not practically adjustable under such conditions. The latter further warrants the accuracy of H<sub>2</sub>O along with CO<sub>2</sub> measurements in cold and/or dry environments.

Accuracy model (10), accuracy Equations 23 and 31, and methodology in Sections 4 and 5 use the composite uncertainty of field CO<sub>2</sub>/H<sub>2</sub>O measurements to define the accuracies for infrared gas analyzers in CPEC systems used in ecosystems. Beyond the applications above, the ultimate objective of this study is to provide an approach for the flux measurement community to specify the accuracy of field CO<sub>2</sub>/H<sub>2</sub>O measurements in ecosystems from the infrared gas analyzers in CPEC systems and eventually in open-path eddy-covariance systems as well.

## Appendix A: Algorithm for Water Vapor Mixing Ratio From Ambient Air Temperature, Relative Humidity, and Atmospheric Pressure

For a given ambient air temperature ( $T_a$  in °C) and atmospheric pressure ( $P$  in kPa), air has a limited capacity to hold a certain amount of water vapor (Wallace & Hobbs, 2006). This capacity is described in terms of saturation water vapor pressure ( $e_s$  in kPa), for moist air, given through the Clausius-Clapeyron equation (Sonntag, 1990):

$$e_s(T_a, P) = \begin{cases} 0.6112 \exp\left(\frac{17.62T_a}{T_a + 243.12}\right) f(P) & T_a \geq 0 \\ 0.6112 \exp\left(\frac{22.46T_a}{T_a + 272.62}\right) f(P) & T_a < 0 \end{cases} \quad (\text{A1})$$

where  $f(P)$  is an enhancement factor for moist air, being a function of atmospheric pressure:  $f(P) = 1.0016 + 3.15 \times 10^{-5}P - 0.0074P^{-1}$ . At relative humidity (RH in %), the water vapor pressure [ $e_{RH}(T_a, P)$ ] is:

$$e_{RH}(T_a, P) = RH e_s(T_a, P) \quad (A2)$$

Given the mole numbers of H<sub>2</sub>O ( $n_{RH}$ ) and dry air ( $n_d$ ) per unit volume at RH, the water vapor mixing ratio at RH ( $\chi_{H_2O}^{RH}$ ):

$$\chi_{H_2O}^{RH} \equiv \frac{n_{RH}}{n_d} = \frac{n_{RH} R^* (T_a + 273.15)}{n_d R^* (T_a + 273.15)} = \frac{e_{RH}(T_a, P)}{P_d} \quad (A3)$$

where  $R^*$  is the universal gas constant and  $P_d$  is dry air pressure. Using this equation and the relation:

$$P = P_d + e_{RH}(T_a, P) \quad (A4)$$

$\chi_{H_2O}^{RH}$  can be expressed in mmolH<sub>2</sub>O mol<sup>-1</sup> as

$$\chi_{H_2O}^{RH} = \frac{1000 \times e_{RH}(T_a, P)}{P - e_{RH}(T_a, P)} \quad (A5)$$

This is used to calculate  $\chi_{H_2O}^{RH}$  in Figures 2 and 4.

## Data Availability Statement

The data that support the findings of this study are openly available at <https://datadryad.org/stash/share/pdcabKdKQdCo9ug-T-oTRC66Lo4uesR2FCT7hO4inT0>.

## Acknowledgments

Authors thank Campbell Scientific for providing the images used in Figure 1, anonymous reviewers for their deep insight into this study topic and constructive comments to improve our presentations, and Linda Worlton-Jones for her professional proofreading. This research was supported by the Strategic Priority Research Program of the Chinese Academy of Sciences (XDA19030204), (2016YFC0500300), Long Term Agroecosystem Network (LTAR-USDA), and Campbell Scientific Research and Development.

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